

Spatially-varying meshless approximation method for enhanced computational efficiency

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Motivation

Realistic problems do not have closed form solutions.

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abla} = 0,$$

 $rac{\partial ec{v}}{\partial t} + ec{v} \cdot
abla ec{v} = -
abla p +
abla \cdot (\operatorname{Ra}
abla ec{v}) - ec{g} \operatorname{Ra} \operatorname{Pr} T_{\Delta},$
 $rac{\partial T}{\partial t} + ec{v} \cdot
abla T =
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abla T)$



Numerical treatment is required:

- 1. Domain discretization
- 2. Differential operator approximation
- 3. PDE discretization
- 4. Solve sparse linear system

Differential operator approximation

$$(\mathcal{L}u)(\boldsymbol{x}_c) \approx \sum_{i=1}^n w_i u(\boldsymbol{x}_i)$$

 $\mathcal{L}\Big|_{\boldsymbol{x}_c} = \boldsymbol{w}_{\mathcal{L}}(\boldsymbol{x}_c)^T$



RBF-FD

- Polyharmonic splines augmented with monomials
- + Higher stability
- + Stable on scattered nodes
- Computationally demanding

Stencil size: 12 in 2D, 20 in 3D

MON

- A set of monomials centred at stencil nodes.
- $+ \ \ Computationally \ cheap$
- Unstable on irregular nodes

Stencil size: 5 in 2D and 7 in 3D

Both allow control over the order of the approximation.

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Spatially-varying approximation method

Spatially-varying approximation method



Computational stability

Computational stability

The stability is estimated via the condition number $\kappa(\mathbf{M}) = \|\mathbf{M}\| \|\mathbf{M}^{-1}\|$, where $\|\cdot\|$ denotes the L^2 norm.



Governing problem

To assess the advantages of the hybrid method, we focus on the <u>natural convection problem</u> that is governed by a system of three PDEs that describe the continuity of mass, the conservation of momentum and the transfer of heat

$$\boldsymbol{\nabla} \cdot \vec{\boldsymbol{v}} = \boldsymbol{0}, \tag{1}$$

$$\frac{\partial \vec{v}}{\partial t} + \vec{v} \cdot \nabla \vec{v} = -\nabla p + \nabla \cdot (\Pr \nabla \vec{v}) - \operatorname{RaPr} \vec{g} T_{\Delta}, \qquad (2)$$

$$\frac{\partial T}{\partial t} + \vec{v} \cdot \nabla T = \nabla \cdot (\nabla T), \tag{3}$$

where a dimensionless nomenclature using Rayleigh (Ra) and Prandtl (Pr) numbers is used.

The de Vahl Davis problem

- To establish confidence in the presented solution procedure.
- Because the regularity of the domain shape allows us to efficiently discretize it using exclusively scattered or regular nodes.



The de Vahl Davis problem: example solution



The de Vahl Davis problem: Nusselt number

Nusselt number: a convenient scalar value for comparison with reference solutions. The average Nusselt number (\overline{Nu}) is calculated as the average of the Nusselt values at the cold wall nodes

$$\mathrm{Nu} = \frac{L}{T_H - T_C} \left| \frac{\partial T}{\partial \boldsymbol{n}} \right|_{\boldsymbol{x}=0}.$$
 (4)



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The de Vahl Davis problem: Comparing with references

Approximation	Nu	execution time [h]	N
scattered	8.867	6.23	55 477
regular	8.852	2.42	64 005
hybrid	8.870	3.11	59 694
Kosec and Šarler (2007)	8.97	/	10201
Sadat and Couturier (2000)	8.828	/	22801
Wan et. al. (2001)	8.8	/	10201

 Table 1: Average Nusselt along the cold edge along with execution times and number of discretization nodes.

• The hybrid method shows significantly shorter computational time (about 50 %), than that required by the scattered discretization employing RBF-FD.

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The de Vahl Davis problem: Convergence analysis



The de Vahl Davis problem: Width of scattered layer δ_h

The domain is split into two parts at a distance $h\delta_h$ from the origin in the lower left corner.

- Horizontal split, resulting in scattered nodes below the imaginary split and regular nodes above it.
- Vertical split, resulting in scattered nodes to the left of it and regular nodes to the right of it.



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Natural convection on irregularly shaped domains



Natural convection on irregularly shaped domains: Nusselt number



The hybrid method effectively reduces the execution time for approximately 35 %. The pure regular discretization with MON approximation is omitted from the table because a stable numerical solution could not be obtained.

Approximation	Nu	execution time [min]	Ν
scattered	12.32	46.31	10 534
hybrid	12.36	29.11	11 535

Table 2: Average Nusselt along the cold duck edges along with execution times. Note that all values in the table were obtained for $\delta_h = 4$.

Application to three-dimensional irregular domains



The scattered method took about 48 hours and the hybrid scattered-regular approximation method took 20 hours to simulate 1 dimensionless time unit with the dimensionless time step $dt = 7.8125 \cdot 10^{-6}$ and about 75 000 computational nodes with $\delta_h = 4$.

Approximation	Nu	execution time [h]	N
scattered	7.36	48.12	65 526
hybrid	6.91	20.54	74 137

Table 3: Average Nusselt along the cold spheres, execution time, and number of computational nodes.

Note: The pure regular discretization with MON approximation is again omitted from the table because a stable numerical solution could not be obtained.

Conclusions

- Proposed a computationally efficient approach to the numerical treatment of PDEs.
- Verified on a solution to a two-dimensional de Vahl Davis natural convection problem.
- Demonstrated on a solution within irregular two- and three-dimensional domains.

Future work:

- * A detailed study of the width of the scattered layer δ_h .
- Further study of aggressiveness of *h*-refinement.
- * Demonstration on more difficult problems: Mixed convection.

Thank you for your attention. Questions?